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1 **Single-junction solar cells based on *p-i-n* GaAsSbN heterostructures grown by liquid**  
2 **phase epitaxy**

3

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18

19 **ABSTRACT**

20 In this paper we present single heterojunction *p-i-n* GaAsSbN/GaAs solar cells grown by low-  
21 temperature liquid-phase epitaxy (LPE) – this is of interest as a component of multi-junction  
22 solar cell devices. The quaternary absorber layer was characterized by low excitation power  
23 photoluminescence to give the temperature dependence of the band gap. This conformed to

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24 the Varshni function at low temperatures to within 10 meV, indicating relatively small alloy  
25 potential fluctuations. The absorption properties and the transport of the photogenerated  
26 carriers in the heterostructures was investigated using surface photovoltage method. A power  
27 conversion efficiency of 4.15 % (AM1.5, 1000 W.m<sup>-2</sup>) was measured for p-i-n  
28 GaAsSbN/GaAs solar cells, which is comparable to the efficiency of MOCVD grown devices  
29 of this type. This is promising for the first report of LPE grown GaAsSbN/GaAs solar cells  
30 since the current record efficiency for the cells based on these compounds grown by MBE  
31 stands just at 6 %. The long-wavelength photosensitivity of the cells determined from external  
32 quantum efficiency and surface photovoltage measurements was shown to be extended to  
33 1040 nm.

34

### 35 **Keywords**

36 GaAsSbN, liquid phase epitaxy, p-i-n heterostructures, solar cells, photovoltaic

37

### 38 **1. Introduction**

39 There has been great interest in dilute nitride III-V-N materials during the last two  
40 decades, driven in part by their potential application in multijunction solar cells (Friedman et  
41 al., 1998; Geisz et al., 2018; Harris, 2005; Isoaho et al., 2019; Johnston et al., 2005; Kurtz et  
42 al., 2002; Miyashita et al., 2013, 2012; Ptak et al., 2009, 2005) which are expected to out-  
43 perform single junction devices. A conversion efficiency of 46.1 % has been reported for  
44 four-junction solar cells under concentrated light, using wafer bonding to combine 2 two-  
45 junction solar cells grown on InP and GaAs substrates (Dimroth et al., 2016). The efficiency  
46 record is currently held by NREL for a six-junction inverted metamorphic concentrator solar cell  
47 which achieved 47.1 % (Geisz et al., 2018). The bandgap combination of the subcells is a key  
48 factor for further improvements in the overall cell efficiency. Presently multi-junction solar

49 cell performance is limited by the performance of the subcells which need to be chosen for  
50 their bandgaps and also need to be grown with appropriate crystal quality. Dilute nitride alloys  
51 such as InGaAsN or GaAsSbN can provide adjustable bandgaps between 1.2 and 0.8 eV  
52 while remaining lattice-matched to GaAs or Ge substrates. Hence the development of these  
53 materials is of great significance for high-efficiency multijunction solar cells, where they can  
54 be used to collect the low-energy photons. However, the device performance has not reached  
55 expectations, due to the low radiative efficiencies and low minority carrier diffusion lengths  
56 (Johnston et al., 2005; Kurtz et al., 2002). The issue of poor minority carrier diffusion length  
57 has been partially tackled by increasing the depletion region width by the use of undoped  
58 layers to increase the current (Miyashita et al., 2012; Ptak et al., 2005). However, the trap-  
59 assisted recombination dark current increases with the width of the depletion region, which  
60 leads to the lowering of the open-circuit voltage.

61 It is therefore essential to improve the properties of dilute nitride compounds. There has  
62 been significant progress in the development of InGaAsN materials. InGaAsN solar cells  
63 grown by metalorganic chemical vapour deposition (MOCVD). A device based on a double-  
64 heterostructure single-junction (and using an anti-reflection coating) has achieved 13.2 %  
65 efficiency as reported by Kim *et al.* (Kim et al., 2015). The Solar Junction Corporation  
66 reported the highest efficiency monolithic triple-junction solar cells of 43.5 % under  
67 concentrated light for using InGaAsN instead of Ge as the lowest subcell (Wiemer et al.,  
68 2011).

69 Another dilute nitride material suitable for solar cell applications as an alternative to  
70 InGaAsN is GaAsSbN. It offers the possibility for independent tuning of the conduction and  
71 valence bands. While the bandgap of dilute nitrides is primarily reduced by lowering the  
72 conduction band minimum, the bandgap of antimonides is reduced by raising the valence  
73 band maximum energy. Both of these mechanisms are explained through the band anti-

74 crossing model. The incorporation of both Sb and N atoms into the crystal lattice enables  
75 lattice matching with GaAs or Ge. In addition, their incorporation also causes a large  
76 concentration of localized states which results in changed electronic and optical properties,  
77 and lower device performance, due to reduced carrier collection. Despite the beneficial  
78 features of this material it has not been as widely studied as InGaAsN, although interest in  
79 GaAsSbN has risen over the last few years (Bian et al., 2004; Gonzalo et al., 2019; T. W. Kim  
80 et al., 2014; Kim, Tae Wan et al., 2014; Lin et al., 2013; Milanova et al., 2019; Tan et al.,  
81 2011; Thomas et al., 2015; Yurong et al., 2017). Recent solar cells based on GaAsSbN have  
82 demonstrated efficiencies of 4 % for a nonoptimized single-junction solar cell structure,  
83 without anti-reflection coating grown by MOCVD ( Kim, Tae Wan et al., 2014) and 6 % for  
84 molecular beam epitaxy (MBE) grown structures (Thomas et al., 2015).

85 In this paper, we present the results of single-heterojunction p-i-n GaAsSbN solar cells  
86 grown by low-cost liquid phase epitaxy (LPE) method. To the best of our knowledge, no other  
87 groups have reported on LPE grown GaAsSbN/GaAs heterostructures. Temperature-  
88 dependent photoluminescence (PL) spectra at a low excitation power of 0.5 mW have been  
89 used to identify the in-gap localized states, which increase the dark current of the cells. In the  
90 first device trials for LPE GaAsSbN/GaAs devices, efficiencies of 4.15 % were achieved  
91 under AM1.5 conditions on areas of 3.5 x 3.5 mm<sup>2</sup>.

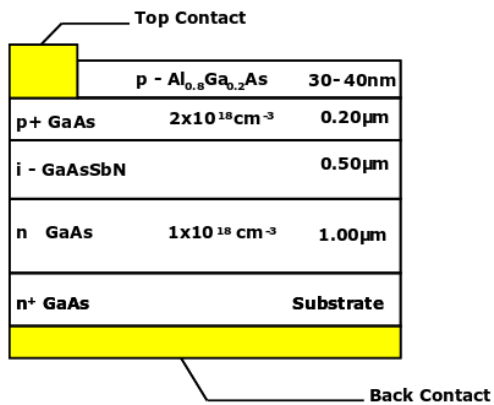
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## 93 **2. Experimental Methods**

94 The schematic structure of the single-junction solar cell investigated in this study is shown in  
95 Fig. 1. It employs a *p-i-n* structure based on compensated GaAsSbN layers grown via LPE.  
96 The epitaxial structures were grown in a horizontal LPE reactor using the “piston boat”  
97 technique. The starting materials were Ga and Sb metal-solvents of 6N purity, and  
98 polycrystalline GaAs and GaN powder used as sources for As and N, respectively. The

99 elements chosen as dopants were Te for n-type doping and Mg for p-type doping. The  
 100 epitaxial structure was grown in the temperature interval 640-546 °C at a cooling rate  
 101 0.8°C/min. Dilute nitride i-GaAsSbN layer was grown from a 7°C supercooled mixed Ga + 5  
 102 at. % Sb solution in the temperature range 558 -550 °C. The GaAs capping and AlGaAs  
 103 “window” layers were grown in the temperature intervals 550-548 °C and 548 -546°C,  
 104 respectively. The Sb content in the GaAs<sub>1-x-y</sub>Sb<sub>x</sub>N<sub>y</sub> layers was measured at several points on a  
 105 cross-section by energy-dispersive X-ray spectroscopy (EDX). The average value of x was  
 106 determined to be 6.7 % and it is confirmed by X-ray photoelectron spectroscopy (XPS)  
 107 measurements. The presence of 0.1% nitrogen in the studied samples was determined by XPS  
 108 measurements using excitation with Mg K $\alpha$  radiation. The lattice mismatch between the LPE  
 109 grown GaAsSbN layers and the GaAs substrate is about 0.48 %, as was determined from X-  
 110 ray diffraction (004)) curves. The details of the structural characterization based on the XRD  
 111 results are presented in Ref. (Milanova et al., 2019)

**Commented [A1]:** The doping concentration in written correctly in the revised figure



113  
 114  
 115

116 **Fig. 1.** Schematic structure of a single-junction *p-i-n* solar cell based on compensated  
117 GaAsSbN heterostructure grown by LPE.

118  
119 Samples were processed into  $3.5 \times 3.5$  mm<sup>2</sup> photovoltaic mesa diodes comprising of a 3  
120 mm diameter optical window using standard mask lithography. Firstly, a common back *n*-  
121 contact was formed via thermal evaporation of an InGe/Au layer, which was subsequently  
122 annealed at 400°C for 1 minute. Following this, top Ohmic contacts were defined via  
123 lithography consisting of 10 µm wide stripes, separated by 90 µm, across the optical window  
124 to ensure homogenous current injection. The contacts were deposited via thermal evaporation,  
125 which consisted of Au/Zn/Au layers, and subsequently followed by a 380°C anneal for 1  
126 minute. Prior to the contact deposition the top Al<sub>0.8</sub>Ga<sub>0.2</sub>As layer was selectively etched using  
127 hydrogen peroxide and citric acid solution. The mesas were then etched using phosphoric  
128 acid, hydrogen peroxide and de-ionized water etchant. They were etched to a sufficient depth  
129 to ensure electrical isolation between adjacent devices.

130 Temperature dependent PL spectra were measured in the temperature range between 10  
131 and 150 K in order to investigate the optical properties of the grown structures. The  
132 excitation was obtained with a laser light having an energy density of 100 nJ/cm<sup>2</sup> (80 MHz  
133 repetition rate and 550 µW laser power) and wavelength of 680 nm provided by  
134 supercontinuum white laser (Fianium WhiteLase) monochromated with a tunable bandpass  
135 transmission filter (Fianium SuperChrome). The light spot diameter on the sample surface  
136 was 130 µm. The sample was mounted in a closed-loop He cryostat and its temperature was  
137 controlled through an Oxford Instruments ITC503 unit. The PL was collected and recorded by  
138 a fibre-coupled spectrometer (BWSpec Glacier X). Surface photovoltage (SPV) spectroscopy  
139 in metal-insulator-semiconductor (MIS) operation mode was undertaken. This technique was

140 applied to study the optical absorption of the structures using the set-up and the measurement  
141 procedure as described elsewhere (Donchev, 2019).

142 J-V measurements were undertaken using a calibrated TS Space Systems solar  
143 simulator with an AM1.5 spectrum at  $1000 \text{ Wm}^{-2}$ . The external quantum efficiency (EQE)  
144 measurements were performed using a Bentham PVE 300 system in the dark, i.e. without  
145 white light bias. A total of 8 different solar cell variants were produced, each sample having  
146 approximately 9 complete cells.

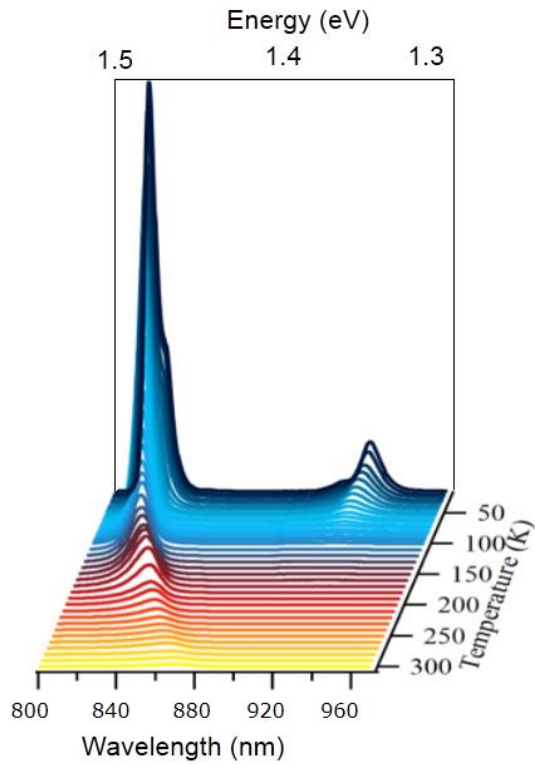
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### 148 **3. Results and Discussion:**

#### 149 *3.1 Photoluminescence Characterization*

150 The temperature-dependent PL spectra of a solar cell epitaxial structure in the range 11  
151 K - 300 K measured under low excitation intensity ( $\sim 0.5 \text{ W/cm}^2$ ) are depicted in Fig. 2. The  
152 peak with high intensity comes from the p+ GaAs layer and the weak peak red-shifted to  
153 GaAs comes from GaAsSbN layer of the structure.





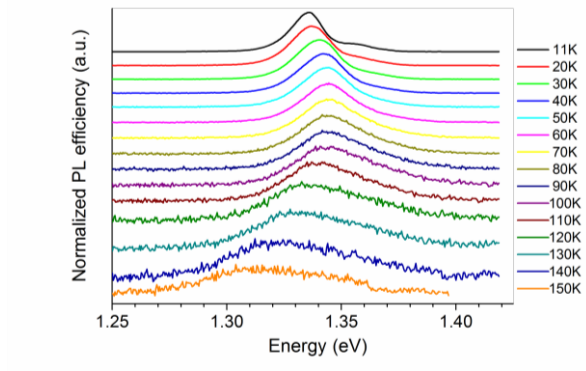
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155 **Fig. 2.** Temperature dependent PL spectra of a GaAsSbN/GaAs solar cell structure in the  
 156 range 15-300 K with step intervals of 10 K.

157

158 Fig. 3 presents the normalized PL peaks of the GaAsSbN layer in the structure  
 159 measured in the range (11 – 150 K). The PL peak energy evolution presents an anomalous,  
 160 nearly S-shaped temperature behaviour. It is seen that there is a blue shift of the PL emission  
 161 energy as the temperature is increased from 10 to about 60-70 K. With further increase of the  
 162 temperature (beyond 70 K) the PL peak exhibits a red shift. This type of behaviour is a well-  
 163 known characteristic of carrier localization effects associated with band-tail states (Gao et al.,  
 164 2016; Lai et al., 2006; Lourenço et al., 2007) which are known to depend on the degree of  
 165 disorder of the compound. In the quaternary GaAsSbN compounds the incorporation of N and

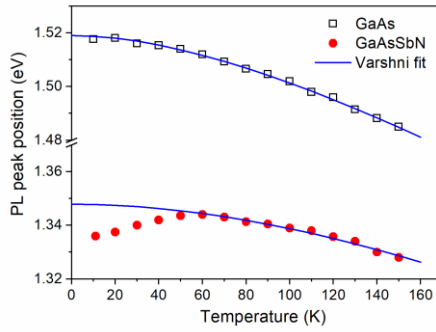
166 Sb into the crystal lattice locally modifies the conduction and valence bands respectively, thus  
 167 creating localized states and potential fluctuations. The blue-shift of the emission at low  
 168 temperatures indicates that as the temperature increases the excitons gain sufficient thermal  
 169 energy to transfer to higher-energy localized levels, thus increasing the emission energy. As  
 170 the temperature increases further (above 70 K) the higher energy localized states are gradually  
 171 saturated, the excitons become almost delocalized and the PL peak energy decreases as a  
 172 function of temperature due to reduction in the bandgap.  
 173



**Fig. 3.** Normalized PL spectra of GaAsSbN in a *p-i-n* solar cell structure measured from 11 K (top) to 150 K (bottom) with step intervals of 10 K under excitation power of 0.5 W/cm<sup>2</sup>.

174  
 175 Fig. 4 illustrates the variations of the PL peak energies as a function of temperature for  
 176 the *p*<sup>+</sup>-GaAs and the GaAsSbN layer in the structure. The PL peak energy for GaAs decreases  
 177 monotonically with increasing temperature. Its temperature dependence is well fitted with  
 178 Varshni's relation  $E_g(T) = E_0 - aT^2/(T+b)$  using the parameters typical for GaAs (Blakemore,  
 179 1982), namely  $E_0 = 1.519$  eV for the bandgap at  $T = 0$  K, and  $a = 5.4 \times 10^{-4}$  eV.K<sup>-1</sup>, and  $b =$   
 180 204 K as fitting parameters.

181



182

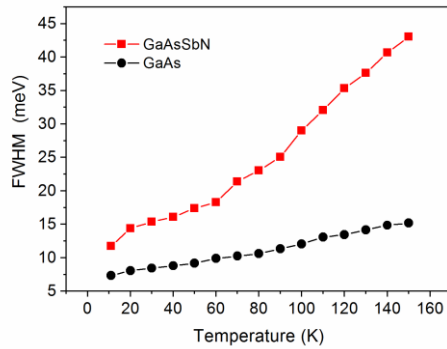
183 **Fig. 4.** Temperature dependence of the PL peak energy of GaAs (squares) and GaAsSbN  
184 (circles) in a *p-i-n* solar cell structure. Lines represent Varshni fits to the data.

185

186 The GaAsSbN peak exhibits the non-typical S-curve (blue-red) behaviour. The first red-  
187 shift is missing (see Fig.4), while the blue shift is only around 10 meV. This indicates that the  
188 potential fluctuations in these samples are relatively small and even at low temperature the  
189 excitons receive enough thermal energy to escape from the localized states and transfer to  
190 higher-energy localized states in the band-tail closer to the conduction band. The fit to the  
191 data above 60 K was obtained using Varshni's relation with the following fitting parameters:  
192  $E_0 = 1.348$  eV,  $a = 6.41 \times 10^{-4}$  eV.K<sup>-1</sup> and  $b = 600$  K. The blue shift of the PL peak energy  
193 observed in this work is compared to the blue shift measured in MBE grown GaAsSbN/GaAs  
194 single quantum wells after annealing. It is nearly the same as the values reported in (Li et al.,  
195 2005) for *in-situ* annealed samples, while in other works (Lourenço et al., 2007; Nunna et al.,  
196 2007) larger values were observed. In our previous work (Milanova et al., 2020) the  
197 temperature-dependent PL spectra of the *p-i-n* structures were measured under higher  
198 excitation intensity (5 W/cm<sup>2</sup>) and no blue shift of the PL peak position was observed at low  
199 temperatures. In the whole temperature range 20 – 300 K, the PL peak energy showed a

200 monotonous decrease with increasing temperature and this behaviour was well fitted by an  
201 empirical Varshni relation (Milanova et al., 2020).

202 Fig. 5 shows the temperature dependence of the PL full widths at half maximum  
203 (FWHM) for  $p^+$  GaAs and GaAsSbN layers. The FWHM values increase monotonically with  
204 temperature from 7.3 meV at 11 K to 15.2 meV at 150 K for GaAs layer. The FWHM for  
205 GaAsSbN is 11.7 meV at 11 K and increases slightly from 14.5 meV to 18 meV in the  
206 temperature range 20 - 60 K where the emission is dominated by localized excitons. FWHM  
207 values increase more rapidly with increasing temperature above 60 K where the emission is  
208 dominated by delocalized excitons.  
209

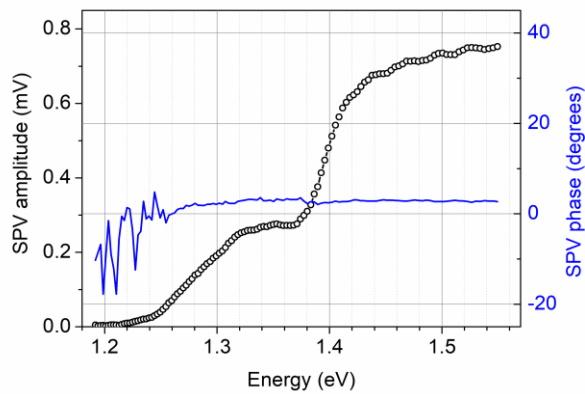


210  
211  
212 **Fig. 5.** Temperature dependence of the FWHM of the PL peak of GaAs (squares) and  
213 GaAsSbN (circles) in a p-i-n solar cell structure.

### 214 3.2. Surface photovoltage characterization

215  
216 SPV spectroscopy has seldom been used to study dilute nitride materials, e.g. the optical  
217 absorption (Bansal et al., 2006) and the band offset (Galluppi et al., 2005) in InGaAsN/GaAs  
218 single quantum wells and the  $E_-$  and  $E_+$  transitions in GaNAs layers (Kudrawiec et al.,

219 2014). However, no other groups have reported on SPV investigations of GaAsSbN dilute  
220 nitride materials. We apply this method to study the optical absorption and photocarrier  
221 transport in the investigated structures. It is well known that in MIS operation mode the SPV  
222 amplitude spectrum emulates the optical absorption spectrum (Kronik and Shapira, 1999),  
223 while the SPV phase spectrum carries information about the direction of the energy band  
224 bending and therefore about the direction of the photocarrier movement (Donchev, 2019). The  
225 SPV measurements were performed at room temperature with a light modulation frequency of  
226 94 Hz. The scanning was from high to low wavelengths, keeping the photon flux constant at  
227 each wavelength.  
228



229  
230  
231 **Fig. 6.** Surface photovoltage amplitude (symbols) and phase (line) spectra of a *p-i-n* solar cell  
232 based on compensated GaAsSbN measured at room temperature.

233  
234 Fig. 6 presents the SPV amplitude and phase spectra of a *p-i-n* single-junction solar  
235 cell structure based on compensated GaAsSbN. The amplitude spectrum reveals a step in the  
236 range of 1.24 – 1.38 eV and another one for energies above 1.38 eV. The former originates

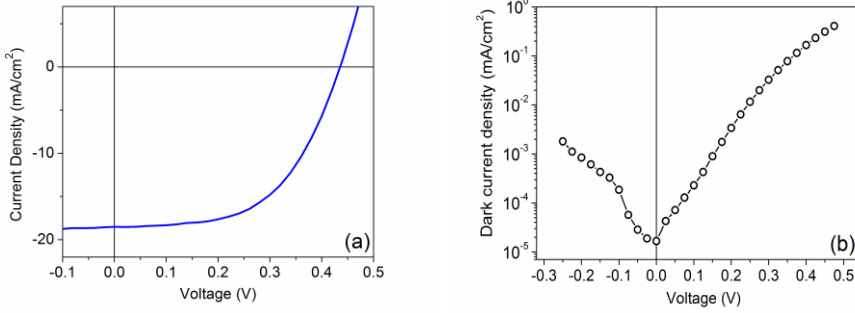
237 from the absorption in the GaAsSbN layer and the latter from absorption in the GaAs layers.  
238 We emphasize that the signals from GaAsSbN and GaAs are comparable in magnitude, which  
239 attests for good quality of the dilute nitride layer. The absorption edge of GaAsSbN from SPV  
240 was confirmed by a Tauc plot as being 1.26 eV. Considering its thickness (0.5  $\mu\text{m}$ ) and Hall  
241 carrier concentration ( $\sim 10^{15} \text{ cm}^{-3}$  (Milanova et al., 2020)) the GaAsSbN layer is fully depleted  
242 and the photogenerated electrons are swept towards the  $n$ -GaAs layer, while the holes –  
243 toward the  $p^+$ -GaAs layer thus giving rise to photovoltage. The direction of the carrier drift is  
244 evidenced by the SPV phase values, which are close to zero degrees in agreement with the  
245 upward energy bands bending (in the direction towards the surface) in the  $p$ - $i$ - $n$  structure  
246 (Donchev, 2019).

247

### 248 3.3 Photovoltaic characterization: J-V and EQE characteristics

249 A typical J-V curve for the GaAsSbN  $p$ - $i$ - $n$  solar cell measured at AM1.5 conditions is  
250 presented in Fig. 7a. Fig. 7b shows the corresponding J-V curve measured in the dark.  
251 Photovoltaic parameters measured on several devices in this work are given in Table 1. The  
252 best solar cell shows an efficiency of 4.15 %, an open-circuit voltage  $V_{oc} = 0.44 \text{ V}$ , short-  
253 circuit current  $J_{sc} = 17.31 \text{ mAcm}^{-2}$  and fill factor  $FF = 54.5 \%$ . The series resistance is  $R_{series} =$   
254  $5.73 \Omega\text{cm}^2$  and shunt resistance is  $R_{shunt} = 478 \Omega\text{cm}^2$ . This is comparable to the efficiency  
255 reported for single-junction 1.25 eV GaAsSbN solar cells with 600 nm thickness of the  
256 GaAsSbN base layer grown by MOCVD (Kim, Tae Wan et al., 2014). A higher efficiency of  
257 about 6 % has been achieved for 1.15 eV MBE grown GaAsSbN solar cells (Thomas et al.,  
258 2015). In both cases, a rapid thermal annealing (RTA) at 800°C of the solar cells was  
259 performed, which significantly increased the open-circuit voltage values to 0.5 – 0.6 V due to  
260 the decrease in the density of the localized states. However, typical values of  $V_{oc}$  measured in  
261 our cells based on as-grown LPE structures without RTA are in the range 0.40 – 0.44 V.

262



263 **Fig. 7.** J-V curves of a GaAsSbN *p-i-n* solar cell measured at AM1.5 conditions (a) and in the  
 264 dark (b).

265

266 Table. 1. PV parameters of several devices measured in this [work](#)

**Commented [A2]:** It does not look right that FF for E407 and E410 are not to 2 decimal points, as is shown for E409

| Sample | Efficiency | $V_{oc}$ [V] | $I_{sc}$ [mA/cm <sup>2</sup> ] | FF [%] |
|--------|------------|--------------|--------------------------------|--------|
| E407   | 4.00       | 0.43         | 17.58                          | 53     |
|        | 4.02       | 0.44         | 17.26                          | 53     |
|        | 4.07       | 0.43         | 17.85                          | 53     |
| E409   | 4.10       | 0.44         | 17.27                          | 53.90  |
|        | 4.09       | 0.44         | 16.92                          | 55.00  |
|        | 4.15       | 0.44         | 17.31                          | 54.50  |
|        | 4.06       | 0.44         | 17.23                          | 53.50  |
| E410   | 4.10       | 0.43         | 18.7                           | 51     |
|        | 4.03       | 0.42         | 18.45                          | 52     |
|        | 3.87       | 0.41         | 18.16                          | 52     |

267

268

269 Theoretically calculated  $V_{oc}$  values for the cells with 1.2 eV bandgap are about 0.8 V

270 assuming that the charge transport is dominated mainly by diffusion of minority carriers.

271 However, even in the best cells, different recombination mechanisms are present which lower

272 the maximum open-circuit voltages below the theoretical limits. It is known that the

273 incorporation even of a small quantity of nitrogen (~0.1%) into GaAs creates a high  
274 concentration of minority carriers traps which sufficiently increases the dark current. To  
275 identify the limiting factors of  $V_{oc}$ , we have determined the saturation dark current from the  
276 dark J-V measurements, which is dominated by trap-assisted recombination in the depletion  
277 region. The saturated dark current density  $J_0$  of the cell obtained from the presented dark J-V  
278 characteristic is around  $10^{-5}$  mA/cm<sup>2</sup>. This value is about three orders of magnitude higher  
279 than those for the high-quality cells which is due to the carrier recombination via N-induced  
280 traps in the GaAsSbN layer. The increase in dark current is directly related to the decreased  
281  $V_{oc}$  values of the solar cells. In addition, the lowering in  $V_{oc}$  values could be explained by a  
282 change in the position of the quasi-Fermi levels from the carrier traps, occurred by the  
283 incorporation of N and Sb into the crystal lattice (Kurtz et al., 2005).  
284

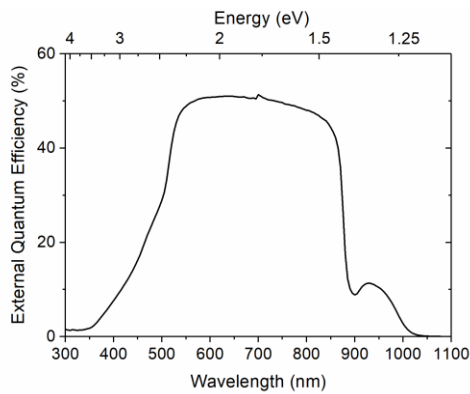


Fig.8. EQE of a GaAsSbN *p-i-n* solar cell.

285  
286 An example of a corresponding EQE graph is given in Fig. 8. EQE values of around 50  
287 % were measured in the wavelength range 550 – 850 nm, which suggests that the upper  $p^+$ -



288 GaAs emitter layer is of good quality despite being grown on GaAsSbN. A significant  
289 decrease of the EQE in the infrared part of the spectrum is due to the short minority carrier  
290 diffusion length in the dilute nitride layer because of the efficient recombination via localized  
291 defect levels. Nevertheless, the EQE extends to approximately 1040 nm (1.19 eV) with an  
292 inflexion point at 1.26 eV in agreement with the SPV results for the bandgap of GaAsSbN. A  
293 slight reduction in EQE is observed at approximately 900 nm (1.377 eV), which can be  
294 associated with the onset of the optical transitions in GaAs in accordance with the SPV  
295 spectrum. The short wavelength photosensitivity of the structure is determined by the  
296 composition and thickness of the AlGaAs layer.

297

#### 298 **4. Conclusions**

299 Single junction solar cells based on a *p-i-n* GaAsSbN/GaAs structure grown via LPE  
300 were developed and studied. Mesa diodes measuring  $3.5 \times 3.5 \text{ mm}^2$  with circular optical  
301 window 3 mm in diameter were fabricated using standard lithography and wet etch  
302 processing. *n*- and *p*- type Ohmic contacts based on InGe/Au and Au/Zn/Au were deposited  
303 via thermal evaporation on the back and the front surface of the cells. Temperature-dependent  
304 PL measurements at low excitation power of 0.5 mW show a slight blue shift of the GaAsSbN  
305 PL emission energy at low temperatures from 10 K to about 70 K, which led to the conclusion  
306 that the potential fluctuations are relatively small. SPV measurements provide information on  
307 the optical absorption and photocarrier transport in the investigated structures. The bandgap  
308 energy at room temperature of GaAsSbN determined from the optical absorption edge is 1.26  
309 eV. Nearly the same IR photosensitivity behaviour was revealed from EQE measurements.

310 J-V curves were measured under standard test conditions (25°C, one sun AM1.5). A  
311 power conversion efficiency of 4.15 %, an open-circuit voltage of 0.44 V, short-circuit  
312 current of 17.31 mA/cm<sup>2</sup> and fill factor 54.5 % were obtained for the cells without anti-

313 reflection coatings and without the rapid thermal annealing that has achieved higher voltages  
314 for MOVPE and MBE materials. This is very promising result for the first LPE grown  
315 GaAsSbN/GaAs solar cells, especially given that these cells are non-optimized, and the record  
316 GaAsSbN/GaAs solar cell currently stands at just 6 % efficiency. Further improvements in  
317 materials quality and in device design are needed to ensure higher photovoltaic performance  
318 of these cells.

319

#### 320 **Declaration of competing interest**

321 The authors declare that they have no known competing financial interests or personal  
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323

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