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Impact of abrupt sea ice loss on Greenland water isotopes during the last glacial period

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Greenland ice cores provide excellent evidence of past abrupt climate changes. However, there is no universally accepted theory of how and why these Dansgaard–Oeschger (DO) events occur. Several mechanisms have been proposed to explain DO events, including sea ice, ice shelf buildup, ice sheets, atmospheric circulation, and meltwater changes. DO event temperature reconstructions depend on the stable water isotope ($\delta^{18}\text{O}$) and nitrogen isotope measurements from Greenland ice cores: interpretation of these measurements holds the key to understanding the nature of DO events. Here, we demonstrate the primary importance of sea ice as a control on Greenland ice core $\delta^{18}\text{O}$: 95% of the variability in $\delta^{18}\text{O}$ in southern Greenland is explained by DO event sea ice changes. Our suite of DO events, simulated using a general circulation model, accurately captures the amplitude of $\delta^{18}\text{O}$ enrichment during the abrupt DO event onsets. Simulated geographical variability is broadly consistent with available ice core evidence. We find an hitherto unknown sensitivity of the $\delta^{18}\text{O}$ paleothermometer to the magnitude of DO event temperature increase: the change in $\delta^{18}\text{O}$ per Kelvin temperature increase reduces with DO event amplitude. We show that this effect is controlled by precipitation seasonality.

abrupt warmings | climate change | Arctic | sea ice | paleoclimate

Sea ice is a key player in the Arctic climate system: it affects precipitation, mass balance, and atmospheric circulation over a large region. Understanding sea ice losses during past abrupt warming events remains challenging (1–7), with the critical relationships between total Arctic (here defined as all Northern Hemisphere) sea ice cover, local climate, and Greenland ice core records still only very poorly understood (8, 9). This is particularly important, since Dansgaard–Oeschger (DO) events are both the largest and best-documented examples of abrupt climate change (10–18).

There has recently been significant progress in reconstructing abrupt DO temperatures increases over Greenland using nitrogen isotopes $\delta^{15}\text{N} - \text{N}_2$ (12, 19). This work indicates jumps in temperature over Greenland of up to 16.5 ± 3 K within a few decades (12, 19). A logical but challenging next step is to elucidate how geographical patterns of change in key ice core records, particularly $\delta^{18}\text{O}$, from Greenland ice cores can be used to provide that crucial missing information on the nature and cause of abrupt warming events, sea ice loss, and its relationship to abrupt temperature rises (20, 21).

DO events are imprinted across the whole of Greenland: wherever last glacial ice is preserved, ice core measurements capture these events (10–12, 19, 22). However, the magnitude of the DO imprint is not identical across the Greenland ice sheet. Early DYE3 ice core measurements suggest that $\delta^{18}\text{O}$ changes during DO warmings may be larger in the south of Greenland (10) compared with central Greenland. More recent ice core data (Fig. 1 *A* and *B*) imply that, while the magnitudes of DO temperature and accumulation changes (from $\delta^{15}\text{N} - \text{N}_2$ and $\delta^{18}\text{O}$) are larger in central Greenland compared with the north and northwest (12), $\delta^{18}\text{O}$ changes could be larger and are likely more variable

in the north and northwest compared with central sites (9, 12, 22). How this spatial variability relates to sea ice loss is currently unknown.

General circulation model (GCM) simulations of $\delta^{18}\text{O}$ enable robust interpretation of records recovered from Greenland ice cores. In particular, they allow us to probe influences on the geographical patterns on the measured $\delta^{18}\text{O}$ change. The ability to decode DO changes from $\delta^{18}\text{O}$ records from Greenland ice cores is thus vital to test ideas about drivers of past abrupt climate change (20, 23–25). Here, we present results from a large ensemble of 32 isotope-enabled GCM simulations of DO-type events.

Our DO-type simulations use a freshwater hosing-type setup. Salt is progressively lost from the North Atlantic during stadial periods; classic hosing (26, 27) mechanisms explain the stadial North Atlantic region cooling that we simulate. After a switch of forcing in the hosing, salt returns to the North Atlantic from the tropical Atlantic Ocean and the wider global ocean. This causes the onset of an abrupt warming DO event.

We generate a suite of stadial climates from a 1,500-y glacial period spin-up simulation and then branch 32 different simulations of DO-type warming events off from this range of stadial climates (Fig. 1 *C–E* and *SI Appendix*, Fig. S1). The simulations feature different amplitudes of effective salt fluxes alongside the range of initial stadial sea ice states. When calculating stadial–interstadial differences, 50 y of data are used for each climate: the 20 y on either side of the salt flux switch are excluded.

Significance

The Dansgaard–Oeschger events contained in Greenland ice cores constitute the archetypal record of abrupt climate change. An accurate understanding of these events hinges on interpretation of Greenland records of oxygen and nitrogen isotopes. We present here the important results from a suite of modeled Dansgaard–Oeschger events. These simulations show that the change in oxygen isotope per degree of warming becomes smaller during larger events. Abrupt reductions in sea ice also emerge as a strong control on ice core oxygen isotopes because of the influence on both the moisture source and the regional temperature increase. This work confirms the significance of sea ice for past abrupt warming events.

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The authors declare no conflict of interest.

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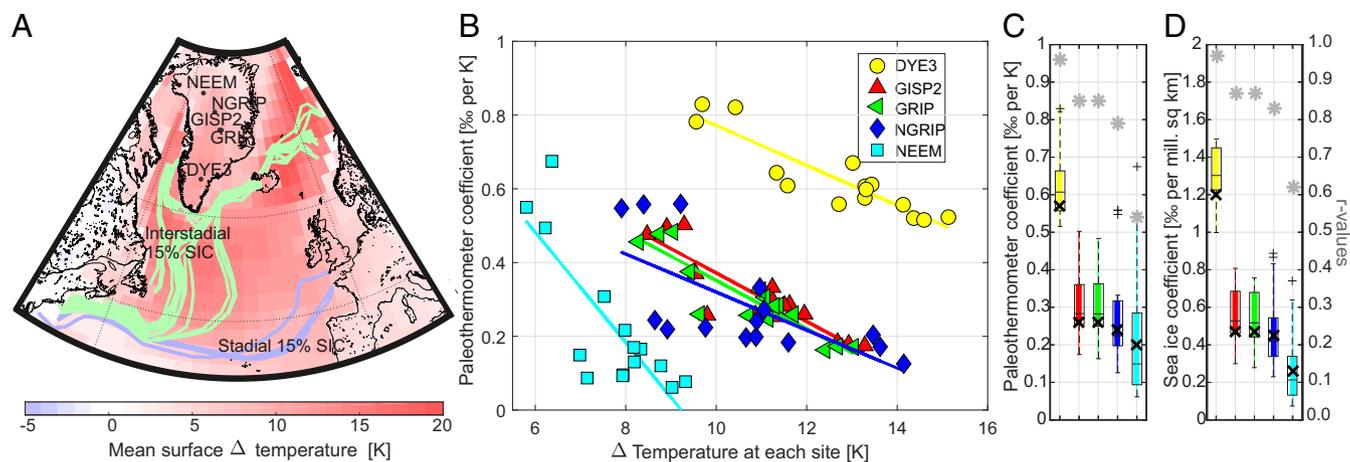


Fig. 2. (A) Mean surface temperature change from simulated stadial to interstadial climates. Averages are from 15 simulations with significant ($+2.0\%$) DO rises in $\delta^{18}\text{O}$ as shown in *SI Appendix, Fig. S1*. Lines show the standard 0.15 mean annual sea ice concentration (SIC) contour for this subset of stadials (blue) and interstadials (green). (B) Paleothermometer values for each individual simulated DO warming event. Coefficients are shown from five ice core sites. Each coefficient is calculated for a single DO event (*Materials and Methods*). Lines indicate the tendency of the paleothermometer values to decrease with the size of warming at each site. Larger variability in these paleothermometer values can be seen at NEEM and to a lesser extent, at NGRIP. (C) The same as B as a boxplot for each ice core site (colors are the same as in B). Any outliers are shown as + symbols. There is a clear increase in the coefficients from the north to the south. (D) The same as C but for sea ice coefficients. Higher coefficients and r values suggest that sea ice reconstructions based on DYE3 ice would be invaluable. Note that r values (gray stars) and coefficients (bold black crosses) derived from least squares best fits shown in *SI Appendix, Fig. S2* are plotted for comparison.

loss (Fig. 3 A–C). A similar geographical pattern also occurs in the temperature change patterns recorded in ice cores (12, 20), and in modeled stadial–interstadial temperature changes (15), lending credence to these simulations. The $\delta^{18}\text{O}$ increase is always strongly positive in the southwest of Greenland, including at DYE3, most commonly with a gradual reduction toward zero change in the northeast. However, some simulated DO events exhibit a stronger east–west geographical gradient in $\delta^{18}\text{O}$ change (Fig. 3B), and others exhibit a stronger south–north gradient (Fig. 3C). Thus, although all simulations have large $\delta^{18}\text{O}$ increases at the most southern DYE3 ice core site, $\delta^{18}\text{O}$ increases are more variable between simulations at the central and northern sites (Fig. 3 A–C). During some simulated events, NEEM changes are larger than those at GISP2 and GRIP and of the same magnitude of those at NGRIP.

Greenland Ice Core Paleothermometers

The $\delta^{18}\text{O}$ –temperature, or traditional paleothermometer (31), coefficient ($\delta^{18}\text{O}$ per Kelvin) over the DO warmings varies considerably across Greenland (Fig. 2 B–D). To reconstruct DO temperature changes from $\delta^{15}\text{N} - \text{N}_2$, a model of firn densification and heat diffusion into the ice is required. This uses initial estimates of accumulation rate, ice age, and temperature, where the latter is derived from $\delta^{18}\text{O}$. Better independent constraints on accumulation rate and temperature from $\delta^{18}\text{O}$ can thus reduce uncertainties on the reconstruction of abrupt DO temperature and accumulation change from $\delta^{15}\text{N} - \text{N}_2$ (12, 19).

It has previously been proposed that the $\delta^{18}\text{O}$ –temperature relationship is dependent on obliquity insolation forcing (19). Here, we investigate another possibility. We calculate the paleothermometer coefficient associated with each individual simulated DO event (i.e., using coefficients from $\Delta\delta^{18}\text{O}$ and Δ temperature) at each ice core site. With this approach applied to 15 simulations with significant ($+2.0\%$) DO rises in $\delta^{18}\text{O}$ (*SI Appendix, Fig. S1*), we obtain a mean paleothermometer coefficient of 0.63 at DYE3; GISP2 yields 0.31. At GRIP, it is 0.30. NGRIP is 0.29, and at NEEM, it is 0.23% per Kelvin. A very similar geographical pattern emerges regardless the type of approach used to calculate the paleothermometer, with high values in the south and low values in the north. Within uncer-

tainties, the simulated coefficients match the few observed coefficients (*Materials and Methods*) (12, 19). Results here support the idea of considerable variability between DO events in the $\delta^{18}\text{O}$ –temperature relationship (12, 19).

Our simulations allow us to decipher what controls $\delta^{18}\text{O}$ –temperature variability between DO events and between ice core sites. In our ensemble, the magnitude of the DO temperature increase has a strong control over the paleothermometer coefficient. Lower $\delta^{18}\text{O}$ –temperature coefficients occur during larger DO events (Fig. 2B) (i.e., we find a strong systematic relationship between the size of the abrupt warming and the paleothermometer coefficient at all Greenland ice core sites). This finding also provides support for the idea that the paleothermometer is fundamentally dependent on the change in temperature at high latitude (32). The pattern varies over Greenland: small decreases occur in the paleothermometer coefficient with warming event size at DYE3 compared with large decreases in the coefficient at NEEM. This identification of a systematic dependence of the $\delta^{18}\text{O}$ –temperature relationship on the size of the abrupt warming may be useful in further constraining Greenland abrupt temperature change records based on $\delta^{15}\text{N} - \text{N}_2$ and $\delta^{18}\text{O}$ measurements.

Precipitation and $\delta^{18}\text{O}$ Seasonality Changes

The current prevailing hypothesis is that most Greenland geographical variability in $\Delta\delta^{18}\text{O}$ is due to geographical variability in the seasonality of precipitation (9, 12, 19). To test this and to better understand the sea ice imprint of DO and the $\delta^{18}\text{O}$ –temperature relationship, we calculate the impact of $\delta^{18}\text{O}$ by only archiving climate information during periods of snow accumulation (33).

If a higher fraction of the yearly precipitation falls during winter months, the $\delta^{18}\text{O}$ record will have a negative $\delta^{18}\text{O}$ bias. In addition to this, the average $\delta^{18}\text{O}$ in each month will also change as the climate moves from a stadial to an interstadial state. We quantify these two effects (by isolating the impact of changes in $\delta^{18}\text{O}$ due to changes in the seasonal cycle of precipitation) and seasonal changes in $\delta^{18}\text{O}$ (34). This is a quantification of local precipitation seasonality ΔP_{seas} vs. other $\Delta\delta^{18}\text{O}_{\text{seas}}$ impacts on changes in $\delta^{18}\text{O}$ ($\Delta\delta^{18}\text{O}$) (*Materials and Methods*).

of a DO-type warming, a large northward North Atlantic flux of salt occurs. This is associated with abruptly increased northward delivery of ocean heat and the melt of substantial areas of North Atlantic sea ice (5, 6). A gradual external forcing related to orographic change of the North American Laurentide ice sheet or a small freshwater perturbation can also trigger similar oscillations (4, 15). A version of the GCM called GFDL CM2Mc also exhibits nonforced oscillations of the AMOC, although only for a certain combination of background climate conditions (39). Here, we use a forced salt (or hosing) oscillation approach for setting up our suite of DO-type simulations.

Simulations are set up using an isotope-enabled version of the Hadley Center Coupled Model general circulation model (HadCM3). This GCM consists of a coupled atmosphere, ocean, and sea ice model and has been widely used to study past, present, and future climates (40, 41). HadCM3 can also be run for hundreds of years on modern supercomputers (24, 42, 43). The ocean component of HadCM3 is a rigid lid model, with a fixed volume and water conservation through a time-invariant surface salinity flux that represents iceberg calving. Ref. 44 has the implementation of water isotope code in HadCM3. Ice sheets and sea ice in the model are initialized with $\delta^{18}\text{O}$ values of -40 and -2‰ , respectively. Using this model, salinity fluxes are applied in opposing directions to the North Atlantic vs. the rest of the global ocean surface. Simulations are set up using LGM ice sheets, orbital forcing, and greenhouse gas composition; additional details are in ref. 42. Every DO simulation is continued from the same standard 1,000-y glacial period (LGM) spin-up simulation. In the first instance, each DO-type simulation is then run for 500 y using an identical stadial-type forcing: a negative salinity flux, equivalent to 0.25 Sv of freshwater, is applied across the North Atlantic between 50° N and 70° N. To balance the global salt (freshwater) budget, a coincident equivalent-sized positive salinity flux is also applied to the rest of the global ocean, ensuring a net global salt flux of zero. This generates a 1,500-y spin up of a stadial-type climate.

A suite of DO simulations is branched from this initial spin-up 1,500 stadial-type simulations. To generate the suite of stadial-type climates, a range of salt fluxes is then applied to the 1,500 spin ups from the equivalent of 0.1- to 1.0-Sv North Atlantic freshwater. The salt flux increments are equivalent to freshwater fluxes of 0.1, 0.25, 0.5, and 1.0 Sv (*SI Appendix, Fig. S1*). Later interstadial negative salinity fluxes are applied using the same range of magnitudes. Each of these stadial-type simulations is branched off from a different year of the spin-up simulation. For each of the salt flux experiments, the surface layer of the ocean is freshened across the North Atlantic at a rate equivalent to the addition of between 0.1 and 1 Sv of freshwater, while the rest of the surface layer of the ocean is salinized at a rate equivalent to the loss of between 0.1 and 1 Sv (i.e., from -0.1 to -1 Sv) of freshwater. The suite of simulations is run using this stadial-type forcing for between 100 and 500 y, yielding a wide range of stadial climates. In each case, the North Atlantic (negative) salt forcing is always balanced by an equivalent-sized (positive) salt forcing applied across the rest of the global ocean; therefore, the net global freshwater (or salt) flux is always zero.

To generate a switch between stadial-type and interstadial-type climate states, a reversed salinity forcing, with positive salinity forcing over the North Atlantic and negative values over the remaining global ocean, is then applied. All forcings are between ± 0.1 and ± 1 Sv. The net global salinity flux is always zero. When we calculate stadial–interstadial difference, we use 50 y of data in each case: the 20 y on either side of the salt flux switch are excluded.

Greenland is represented by 80 grid points in HadCM3; thus, some smoothing of the surface topography is required to run the simulation (43). The modeled surface elevation is thus lower than the observed elevation. For the northern ice core sites, our modeled surface elevations are generally within 500 m of the present day surface elevations: NEEM: 2,450 vs. 2,341 m above sea level (observed vs. modeled, respectively). Similarly, NGRIP is 2,917 vs. 2,788 m above sea level (observed vs. modeled, respectively), and GISP2 is 3,216 vs. 2,832 m above sea level (observed vs. modeled, respectively). Because Greenland is somewhat narrower in the south, the elevations at DYE3 are 2,480 vs. 1,240 m for observed vs. modeled surface elevation, respectively. Thus, while all modeled sites are somewhat too low, the discrepancy is most significant in the south at DYE3. This may account for some, but not all, of the larger sensitivity of $\delta^{18}\text{O}$ at DYE3 during DO events.

The suite of simulations shows a range of $\delta^{18}\text{O}$ values and sea ice states (Fig. 2A and *SI Appendix, Fig. S1*). At NGRIP, the stadial climate $\delta^{18}\text{O}$ varies from -37.5 to -29‰ . Of 32 initial simulations, 15 exhibit DO-type abrupt warming over Greenland with $\delta^{18}\text{O}$ increases of 2.0‰ or more, where a 2.0‰ threshold could be considered representative of the minimum for a

small DO-type abrupt warming (45). Each simulation with a significant $\delta^{18}\text{O}$ jump starts from $\delta^{18}\text{O}$ values at NGRIP of -32‰ or below (*SI Appendix, Fig. S1*). Most of the larger $\delta^{18}\text{O}$ jumps occur under an interstadial-type salt flux (i.e., a return of salt to the North Atlantic) of size equivalent to 0.25 Sv or larger. This magnitude of salt oscillation is also in agreement with the size of salt flux required to induce significant AMOC variations in HadCM3 under present and future forcing scenarios (37). Note that there is a mean offset in Greenland between the mean model and mean data in $\delta^{18}\text{O}$ of $+8\text{‰}$ (Fig. 1 C–E). This offset is in the same direction as in earlier isotopic simulations alongside a similar warm model bias in the temperature and a wet bias in the precipitation (21).

Temperature and Sea Ice Coefficients. There are two main approaches that are used to calculate $\delta^{18}\text{O}$ –temperature coefficients and similarly, $\delta^{18}\text{O}$ –sea ice coefficients. Note that the terms gradients and slopes are also used to denote these coefficient values. The first and most common approach to calculating the coefficients is to fit a second-degree polynomial to a set of stadial and interstadial $\delta^{18}\text{O}$ and temperature values, minimizing the least squares term (46, 47). Coefficients (or gradients or slopes) from this method are shown in Fig. 2 C and D (bold black crosses) and *SI Appendix, Figs. S2 and S3*.

The second approach is to calculate the coefficients for each individual simulated DO event so that the coefficient in this case is simply $\Delta\delta^{18}\text{O}/\Delta$ temperature for the paleothermometer. Additionally, $\Delta\delta^{18}\text{O}/\Delta$ Arctic sea ice area indicates the sea ice coefficient. Note that Arctic sea ice area is a total Northern Hemisphere value that is calculated by summing each grid box multiplied by the sea ice concentration and grid box area. This approach yields a distribution of coefficients and is arguably a more accurate depiction of the relationship between $\Delta\delta^{18}\text{O}$, Δ temperature, and Δ Arctic sea ice area for the suite of simulated DO events. This approach is used to characterize the distribution of paleothermometer and sea ice coefficients (Fig. 2 B–D, colored and boxplot results).

Using the second approach and the same set of the largest DO events, we obtain a mean paleothermometer coefficient of 0.63‰ per Kelvin (16th to 84th percentile is 0.53 to 0.77) at DYE3; GISP2 yields 0.31‰ per Kelvin (0.20 to 0.46), GRIP is 0.30‰ per Kelvin (0.19 to 0.45), NGRIP is 0.29‰ per Kelvin (0.19 to 0.52), and NEEM is 0.23‰ per Kelvin (0.08 to 0.52). If the first approach is used, a similar geographical pattern and similar values emerge (Fig. 2C). This same pattern emerges independent of whether all or subsets of simulations are used in the calculations, although eliminating the smaller DO events does tend to raise the average values of coefficients. While the average simulated NGRIP coefficient is lower than the overall coefficient of 0.52‰ per Kelvin suggested for NGRIP (19), other measurements imply a somewhat lower coefficient (12), and our set of simulations indicates that the observed coefficient of 0.52 at NGRIP occurs around the 84th percentile value (Fig. 2C) (i.e., within the central range of values of simulated coefficients). Our range also encompasses the interevent set of approximate initial NGRIP coefficients from 0.28 to 0.42‰ per Kelvin used by ref. 19.

Uncertainties on Table Values. The data in Table 1 are mean values from the subset of simulations that have significant (2.0‰) DO increases in $\delta^{18}\text{O}$ at NGRIP. The uncertainties are ± 1 SD calculated from this sample size of 15 DO events.

Isolating the Impact of Changes in Seasonality. To qualify the relative impact of the precipitation vs. $\delta^{18}\text{O}$ seasonal changes, we isolate the impact of changes in $\delta^{18}\text{O}$ due to changes in the seasonal cycle of precipitation and seasonal changes in $\delta^{18}\text{O}$ (34) (Fig. 3 D–F and *SI Appendix, Fig. S4*):

$$\Delta P_{\text{seas}} = \frac{\sum_j \delta^{18}\text{O}_j^{\text{stadial}} \cdot P_j}{\sum_j P_j} - \frac{\sum_j \delta^{18}\text{O}_j^{\text{stadial}} \cdot P_j^{\text{stadial}}}{\sum_j P_j^{\text{stadial}}} \quad [1]$$

Changes in $\delta^{18}\text{O}$ due solely to monthly changes in $\delta^{18}\text{O}$ (nonprecipitation seasonality influences) are calculated:

$$\Delta \delta^{18}\text{O}_{\text{seas}} = \frac{\sum_j \delta^{18}\text{O}_j \cdot P_j^{\text{stadial}}}{\sum_j P_j^{\text{stadial}}} - \frac{\sum_j \delta^{18}\text{O}_j^{\text{stadial}} \cdot P_j^{\text{stadial}}}{\sum_j P_j^{\text{stadial}}} \quad [2]$$

where the summations are over 12 mo (index j). The superscript *stadial* indicates values from the cool preceding stadial-type simulated climate, and no superscript indicates values from the postwarming interstadial-type climate. This method is somewhat different from that used to decompose $\delta^{18}\text{O}$ changes in ref. 48, where daily outputs were used. In particular, residuals

[i.e., $R_{\text{residual}} = (\Delta P_{\text{seas}} + \Delta \delta^{18}\text{O}_{\text{seas}}) - \Delta \delta^{18}\text{O}$] could indicate changes in the higher (than monthly)-frequency covariance between precipitation and $\delta^{18}\text{O}$ not captured by this monthly “seasonal”-type decomposition. However, the R_{residual} is smaller than $\pm 2\%$ over Greenland (less than $\pm 0.5\%$ over most of central Greenland). The stadial to interstadial values are calculated using 50 y of data from each climate, with the 20 y on either side of the abrupt warming excluded from the analysis.

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